

$\lim_{\ell \rightarrow \infty} (\text{AdS}_3/\text{CFT}_2)$   
Flat Space Holography

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TU Wien

All about AdS3  
ETH Zurich, November 2015



## Some of our papers on flat space holography



A. Bagchi, D. Grumiller and W. Merbis,  
“Stress tensor correlators in three-dimensional gravity,”  
arXiv:1507.05620.



A. Bagchi, R. Basu, D. Grumiller and M. Riegler,  
“Entanglement entropy in Galilean conformal field theories and flat  
holography,”  
Phys. Rev. Lett. **114** (2015) 11, 111602 [arXiv:1410.4089].



H. Afshar, A. Bagchi, R. Fareghbal, D. Grumiller and J. Rosseel,  
“Spin-3 Gravity in Three-Dimensional Flat Space,”  
Phys. Rev. Lett. **111** (2013) 12, 121603 [arXiv:1307.4768].



A. Bagchi, S. Detournay, D. Grumiller and J. Simon,  
“Cosmic Evolution from Phase Transition of Three-Dimensional Flat Space,”  
Phys. Rev. Lett. **111** (2013) 18, 181301 [arXiv:1305.2919].



A. Bagchi, S. Detournay and D. Grumiller,  
“Flat-Space Chiral Gravity,”  
Phys. Rev. Lett. **109** (2012) 151301 [arXiv:1208.1658].

# Outline

Motivations

Flat space holography basics

Recent results

Generalizations & open issues

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## Motivations

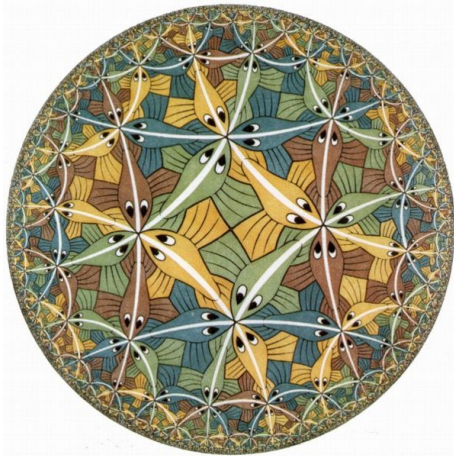
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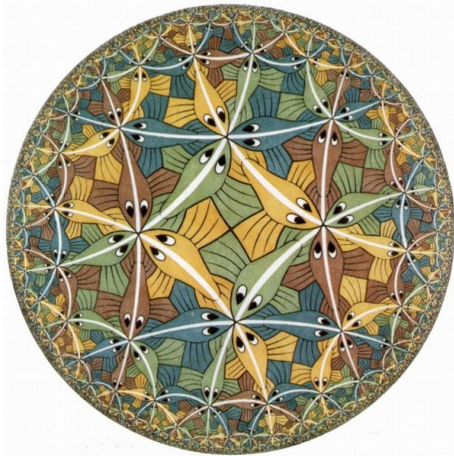
## Holography beyond AdS/CFT?

This talk focuses on holography (in the quantum gravity sense).



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Main question: how general is holography?

### How general is holography?

- ▶ To what extent do (previous) lessons rely on the particular constructions used to date?
- ▶ Are they tied to stringy effects and to string theory in particular, or are they general lessons for quantum gravity?

see numerous talks at KITP workshop “Bits, Branes, Black Holes” 2012

and at ESI workshop “Higher Spin Gravity” 2012

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- ▶ originally holography motivated by unitarity
- ▶ plausible AdS/CFT-like correspondence could work non-unitarily
- ▶ AdS<sub>3</sub>/log CFT<sub>2</sub> first example of non-unitary holography DG, (Jackiw), Johansson '08; Skenderis, Taylor, van Rees '09; Henneaux, Martinez, Troncoso '09; Maloney, Song, Strominger '09; DG, Sachs/Hohm '09; Gaberdiel, DG, Vassilevich '10; ... DG, Riedler, Rosseel, Zojer '13
- ▶ recent proposal by Vafa '14

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- ▶ Can we establish a flat space holographic dictionary?

the answer appears to be yes — see my current talk and recent papers by Bagchi et al., Barnich et al., Strominger et al., '12-'15

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- ▶ Generic non-AdS holography/higher spin holography?

non-trivial hints that it might work

Gary, DG Rashkov '12; Afshar et al '12; Gutperle et al '14-'15; Gary, DG, Prohazka, Rey '14; Lei, Ross '15; Lei, Peng '15; Breunholder, Gary, DG, Prohazka '15; ...

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- ▶ Address questions above in simple class of 3D toy models
- ▶ Exploit gauge theoretic Chern–Simons formulation
- ▶ Restrict to kinematic questions, like (asymptotic) symmetries

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if holography is true  $\Rightarrow$  must work in flat space

Just take large AdS radius limit of  $10^4$  AdS/CFT papers?

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- ▶ Take linear combinations of Virasoro generators  $\mathcal{L}_n, \bar{\mathcal{L}}_n$

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$$[L_n, L_m] = (n - m) L_{n+m} + \frac{c_L}{12} (n^3 - n) \delta_{n+m, 0}$$

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- ▶ This is nothing but the  $BMS_3$  algebra (or  $GCA_2$ ,  $URCA_2$ ,  $CCA_2$ )!

Ashtekar, Bicak, Schmidt, '96; Barnich, Compere '06

$L_n$ : diffeos of circle,  $M_n$ : supertranslations,  $c_{L/M}$ : central extensions

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If dual field theory exists it must be a 2D Galilean CFT!

Bagchi et al., Barnich et al.

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- ▶ Example where it does not work: highest weight conditions

## Flat space Einstein gravity as $isl(2)$ Chern–Simons theory

For details, references and spin-3 generalization see [Gary, DG, Riegler, Rosseel '14](#)

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Achúcarro, Townsend '86; Witten '88



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- ▶ Flat space:  $\mathfrak{isl}(2)$  gauge algebra

$$S_{\text{CS}}^{\text{flat}} = \frac{k}{4\pi} \int \langle \mathcal{A} \wedge d\mathcal{A} + \frac{2}{3} \mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} \rangle$$

with  $\mathfrak{isl}(2)$  connection ( $a = 0, \pm 1$ )

$$\mathcal{A} = e^a M_a + \omega^a L_a$$

$\mathfrak{isl}(2)$  algebra (global part of BMS/GCA)

$$[L_a, L_b] = (a - b)L_{a+b}$$

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Note:  $e^a$  dreibein,  $\omega^a$  (dualized) spin-connection

Bulk EOM: gauge flatness  $\rightarrow$  Einstein equations

$$\mathcal{F} = d\mathcal{A} + \mathcal{A} \wedge \mathcal{A} = 0$$

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- ▶ metric

$$g_{\mu\nu} \sim \frac{1}{2} \text{tr} \langle \mathcal{A}_\mu \mathcal{A}_\nu \rangle \quad \rightarrow \quad ds^2 = M du^2 - 2 du dr + 2N du d\varphi + r^2 d\varphi^2$$

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## Correlation functions in flat space holography

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Start slowly with 0-point function

## 0-point function (on-shell action)

Not check of flat space holography but interesting in its own right

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$$\Gamma = -\frac{1}{16\pi G_N} \int d^3x \sqrt{g} R - \frac{1}{8\pi G_N} \int d^2x \sqrt{\gamma} K - I_{\text{counter-term}}$$

with  $I_{\text{counter-term}}$  chosen such that

$$\delta\Gamma|_{\text{EOM}} = 0$$

for all  $\delta g$  that preserve flat space bc's

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Result (Detournay, DG, Schöller, Simon '14):

$$\Gamma = -\frac{1}{16\pi G_N} \int d^3x \sqrt{g} R - \underbrace{\frac{1}{16\pi G_N}}_{\frac{1}{2} \text{GHY!}} \int d^2x \sqrt{\gamma} K$$

follows also as limit from AdS using Mora, Olea, Troncoso, Zanelli '04  
independently confirmed by Barnich, Gonzalez, Maloney, Oblak '15

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Standard procedure (Gibbons, Hawking '77; Hawking, Page '83)

Evaluate Euclidean partition function in semi-classical limit

$$Z(T, \Omega) = \int \mathcal{D}g e^{-\Gamma[g]} = \sum_{g_c} e^{-\Gamma[g_c(T, \Omega)]} \times Z_{\text{fluct.}}$$

path integral bc's specified by temperature  $T$  and angular velocity  $\Omega$

Two Euclidean saddle points in same ensemble if

- ▶ same temperature  $T = 1/\beta$  and angular velocity  $\Omega$
- ▶ obey flat space boundary conditions
- ▶ solutions without conical singularities

Periodicities fixed:

$$(\tau_E, \varphi) \sim (\tau_E + \beta, \varphi + \beta\Omega) \sim (\tau_E, \varphi + 2\pi)$$

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3D Euclidean Einstein gravity: for each  $T, \Omega$  two saddle points:

- ▶ Hot flat space

$$ds^2 = d\tau_E^2 + dr^2 + r^2 d\varphi^2$$

- ▶ Flat space cosmology

$$ds^2 = r_+^2 \left(1 - \frac{r_0^2}{r^2}\right) d\tau_E^2 + \frac{r^2 dr^2}{r_+^2 (r^2 - r_0^2)} + r^2 \left(d\varphi - \frac{r_+ r_0}{r^2} d\tau_E\right)^2$$

shifted-boost orbifold, see [Cornalba, Costa '02](#)



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- ▶ Calculate the **full** on-shell action  $\Gamma$
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- ▶ Plug two Euclidean saddles in on-shell action and compare free energies

$$F_{\text{HFS}} = -\frac{1}{8G_N} \quad F_{\text{FSC}} = -\frac{r_+}{8G_N}$$

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- ▶ Result of this comparison
  - ▶  $r_+ > 1$ : FSC dominant saddle
  - ▶  $r_+ < 1$ : HFS dominant saddle

Critical temperature:

$$T_c = \frac{1}{2\pi r_0} = \frac{\Omega}{2\pi}$$

HFS “melts” into FSC at  $T > T_c$

Bagchi, Detournay, DG, Simon '13

## 1-point functions (conserved charges)

First check of entries in holographic dictionary: identification of sources and vevs

In  $\text{AdS}_3$ :

$$\delta\Gamma|_{\text{EOM}} \sim \int_{\partial\mathcal{M}} \text{vev} \times \delta \text{source} \sim \int_{\partial\mathcal{M}} T_{\text{BY}}^{\mu\nu} \times \delta g_{\mu\nu}^{\text{NN}}$$

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- ▶ analogue of Brown–York stress tensor?
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- ▶ comparison with canonical results

everything works (Detournay, DG, Schöller, Simon, '14)

mass and angular momentum:

$$M = \frac{g_{tt}}{8G} \quad N = \frac{g_{t\varphi}}{4G}$$

full tower of canonical charges: see Barnich, Compere '06

## 2-point functions (anomalous terms)

First check sensitive to central charges in symmetry algebra

Galilean CFT on cylinder ( $\varphi \sim \varphi + 2\pi$ ):

$$\langle M(u_1, \varphi_1) M(u_2, \varphi_2) \rangle = 0$$

$$\langle M(u_1, \varphi_1) N(u_2, \varphi_2) \rangle = \frac{c_M}{2s_{12}^4}$$

$$\langle N(u_1, \varphi_1) N(u_2, \varphi_2) \rangle = \frac{c_L - 2c_M\tau_{12}}{2s_{12}^4}$$

with  $s_{ij} = 2 \sin[(\varphi_i - \varphi_j)/2]$ ,  $\tau_{ij} = (u_i - u_j) \cot[(\varphi_i - \varphi_j)/2]$

Fourier modes of Galilean CFT stress tensor on cylinder:

$$M := \sum_n M_n e^{-in\varphi} - \frac{c_M}{24}$$

$$N := \sum_n (L_n - inuM_n) e^{-in\varphi} - \frac{c_L}{24}$$

Conservation equations:  $\partial_u M = 0$ ,  $\partial_u N = \partial_\varphi M$

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Summarize first how this works in the AdS case

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Illustrate shortcut in  $\text{AdS}_3/\text{CFT}_2$  (restrict to one holomorphic sector)

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Drinfeld, Sokolov '84, Polyakov '87, H. Verlinde '90  
Bañados, Caro '04

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Apply shortcut to flat space/Galilean CFT (Bagchi, DG, Merbis '15)

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- ▶ Correct 2-point functions for Einstein gravity with  $c_L = 0$ ,  $c_M = 12k$

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- ▶ Result on gravity side matches precisely Galilean CFT results

$$\langle M^1 N^2 N^3 \rangle = \frac{c_M}{s_{12}^2 s_{13}^2 s_{23}^2} \quad \langle N^1 N^2 N^3 \rangle = \frac{c_L - c_M \tau_{123}}{s_{12}^2 s_{13}^2 s_{23}^2}$$

provided we choose again the Einstein values  $c_L = 0$  and  $c_M = 12k$

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First correlators with non-universal function of cross-ratios

- ▶ Repeat this algorithm, localizing the sources at three points

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- ▶ Derive 4-point functions for Galilean CFTs (Bagchi, DG, Merbis '15)

$$\langle M^1 N^2 N^3 N^4 \rangle = \frac{2c_M g_4(\gamma)}{s_{14}^2 s_{23}^2 s_{12} s_{13} s_{24} s_{34}}$$

$$\langle N^1 N^2 N^3 N^4 \rangle = \frac{2c_L g_4(\gamma) + c_M \Delta_4}{s_{14}^2 s_{23}^2 s_{12} s_{13} s_{24} s_{34}}$$

with the cross-ratio function

$$g_4(\gamma) = \frac{\gamma^2 - \gamma + 1}{\gamma} \quad \gamma = \frac{s_{12} s_{34}}{s_{13} s_{24}}$$

and

$$\Delta_4 = 4g_4'(\gamma)\eta_{1234} - (\tau_{1234} + \tau_{14} + \tau_{23})g_4(\gamma)$$
$$\eta_{1234} = \sum (-1)^{1+i-j} (u_i - u_j) \sin(\varphi_k - \varphi_l) / (s_{13}^2 s_{24}^2)$$

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$$\langle M^1 N^2 N^3 N^4 N^5 \rangle = \frac{4c_M g_5(\gamma, \zeta)}{\prod_{1 \leq i < j \leq 5} s_{ij}}$$
$$\langle N^1 N^2 N^3 N^4 N^5 \rangle = \frac{4c_L g_5(\gamma, \zeta) + c_M \Delta_5}{\prod_{1 \leq i < j \leq 5} s_{ij}}$$

with the previous definitions and ( $\zeta = \frac{s_{25} s_{34}}{s_{35} s_{24}}$ )

$$g_5(\gamma, \zeta) = \frac{\gamma + \zeta}{2(\gamma - \zeta)} - \frac{(\gamma^2 - \gamma\zeta + \zeta^2)}{\gamma(\gamma - 1)\zeta(\zeta - 1)(\gamma - \zeta)} \times ([\gamma(\gamma - 1) + 1][\zeta(\zeta - 1) + 1] - \gamma\zeta)$$

$$\Delta_5 = 4\partial_\gamma g_5(\gamma, \zeta) \eta_{1234} + 4\partial_\zeta g_5(\gamma, \zeta) \eta_{2345} - 2g_5(\gamma, \zeta) \tau_{12345}$$

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## $n$ -point functions (holographic Ward identities and recursion relations)

Shortcut to 42 (Bagchi, DG, Merbis '15)

Smart check of all  $n$ -point functions?

- ▶ Idea: calculate  $n$ -point function from  $(n - 1)$ -point function

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- ▶ After small derivation we find ( $c_{ij} := \cot[(\varphi_i - \varphi_j)/2]$ )

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## $n$ -point functions (holographic Ward identities and recursion relations)

Shortcut to 42 (Bagchi, DG, Merbis '15)

Smart check of all  $n$ -point functions?

- ▶ Idea: calculate  $n$ -point function from  $(n - 1)$ -point function
- ▶ Need Galilean CFT analogue of BPZ-recursion relation

$$\langle T^1 T^2 \dots T^n \rangle = \sum_{i=2}^n \left( \frac{2}{s_{1i}^2} + \frac{c_{1i}}{2} \partial_{\varphi_i} \right) \langle T^2 \dots T^n \rangle + \text{disconnected}$$

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- ▶ We can also derive same recursion relations on gravity side!

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Fairly non-trivial check that 3D flat space holography can work!

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Works! (Bagchi, Detournay, Fareghbal, Simon '13, Barnich '13)

$$S_{\text{gravity}} = S_{\text{BH}} = \frac{\text{Area}}{4G_N} = 2\pi h_L \sqrt{\frac{c_M}{2h_M}} = S_{\text{GCFT}}$$

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$$S_{\text{EE}}^{\text{GCFT}} = \underbrace{\frac{c_L}{6} \ln \frac{\ell_x}{a}}_{\text{like CFT}} + \underbrace{\frac{c_M}{6} \frac{\ell_y}{\ell_x}}_{\text{like grav anomaly}}$$

Calculation on gravity side confirms result above  
(using Wilson lines in CS formulation)

# Outline

Motivations

Flat space holography basics

Recent results

Generalizations & open issues

## Generalizations & open issues

Recent generalizations:

- ▶ adding chemical potentials

Works! (Gary, DG, Riegler, Rosseel '14)

In CS formulation:

$$A_0 \rightarrow A_0 + \mu$$

Recent generalizations:

- ▶ adding chemical potentials
- ▶ 3-derivative theory: flat space chiral gravity (Bagchi, Detournay, DG '12)

Conformal CS gravity at level  $k = 1$  with flat space boundary conditions conjectured to be dual to chiral half of monster CFT.

Action (gravity side):

$$I_{\text{CSG}} = \frac{k}{4\pi} \int d^3x \sqrt{-g} \varepsilon^{\lambda\mu\nu} \Gamma^\rho{}_{\lambda\sigma} (\partial_\mu \Gamma^\sigma{}_{\nu\rho} + \frac{2}{3} \Gamma^\sigma{}_{\mu\tau} \Gamma^\tau{}_{\nu\rho})$$

Partition function (field theory side, see Witten '07):

$$Z(q) = J(q) = \frac{1}{q} + 196884 q + \mathcal{O}(q^2)$$

Note:  $\ln 196883 \approx 12.2 = 4\pi + \text{quantum corrections}$

## Generalizations & open issues

Recent generalizations:

- ▶ adding chemical potentials
- ▶ 3-derivative theory: flat space chiral gravity (Bagchi, Detournay, DG '12)
- ▶ generalization to supergravity

Works! (Barnich, Donnay, Matulich, Troncoso '14)

Asymptotic symmetry algebra = super-BMS<sub>3</sub>

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Remarkably it exists! (Afshar, Bagchi, Fareghbal, DG, Rosseel '13; Gonzalez, Matulich, Pino, Troncoso '13)

New type of algebra: W-like BMS (“BMW”)

$$[U_n, U_m] = (n - m)(2n^2 + 2m^2 - nm - 8)L_{n+m} + \frac{192}{c_M}(n - m)\Lambda_{n+m} \\ - \frac{96(c_L + \frac{44}{5})}{c_M^2}(n - m)\Theta_{n+m} + \frac{c_L}{12}n(n^2 - 1)(n^2 - 4)\delta_{n+m,0}$$

$$[U_n, V_m] = (n - m)(2n^2 + 2m^2 - nm - 8)M_{n+m} + \frac{96}{c_M}(n - m)\Theta_{n+m} \\ + \frac{c_M}{12}n(n^2 - 1)(n^2 - 4)\delta_{n+m,0}$$

$$[L, L], [L, M], [M, M] \text{ as in BMS}_3 \quad [L, U], [L, V], [M, U], [M, V] \text{ as in isl}(3)$$



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Barnich, Gonzalez, Maloney, Oblak '15: 1-loop partition function matches  $BMS_3$  character

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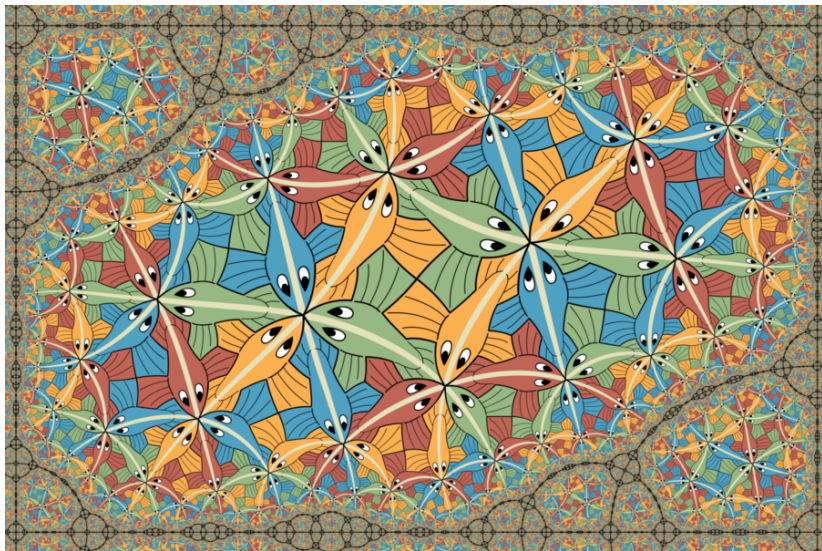
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- ▶ holography seems to work in flat space
- ▶ holography more general than AdS/CFT
- ▶ (when) does it work even more generally?

Thanks for your attention!



Vladimir Bulatov, M.C. Escher Circle Limit III in a rectangle